

## Numerical Solution of Partial Differential Equations:

*The general second – order linear partial differential equation is of the form:*

$$A \frac{\partial^2 u}{\partial x^2} + B \frac{\partial^2 u}{\partial x \partial y} + C \frac{\partial^2 u}{\partial y^2} + D \frac{\partial u}{\partial x} + E \frac{\partial u}{\partial y} + Fu = G$$

*Which can be written as :*

$$Au_{xx} + Bu_{xy} + Cu_{yy} + Du_x + Eu_y + Fu = G \quad \text{.....(1)}$$

*Where A,B,C,.....,G are all functions of x & y*

$$au_{xx} + 2bu_{xy} + cu_{yy} = F(x, y, u, u_x, u_y)$$

*u is the unknown function*

*Equations of form (1) can be classified with respect to the sign of the discriminant:*

$$\Delta s = B^2 - 4AC$$

*in the following ways:*

- (1) If  $\Delta s < 0$  at a point in the  $(x, y)$  plane, the equation is said to be elliptic type.*
- (2) If  $\Delta s > 0$  at that point is said to be hyperbolic type.*
- (3) Parabolic type when  $\Delta s = 0$ .*

*Elliptic type*       $4ac - b^2 > 0$       *Laplace equation*

*Parabolic type*       $4ac - b^2 = 0$       *Heat equation*

*Hyperbolic type*       $4ac - b^2 < 0$       *Wave equation*

*In the following, we will restrict our solves to three simple particular cases of Eq. (1); namely:*

$$u_{xx} + u_{yy} = 0 \quad (\text{the Laplace equation})$$

$$u_{xx} - \frac{1}{c^2} u_{tt} = 0 \quad (\text{the wave equation})$$

$$u_{xx} - u_t = 0 \quad (\text{the heat conduction equation})$$

## Finite – Difference Approximations to derivatives:

Let the  $(x, y)$  plane be divided into a network of rectangles of sides  $\underline{\Delta x = h}$  and  $\underline{\Delta y = k}$  by drawing the sets of lines:

$$x = ih; \quad i = 0, 1, 2, 3, \dots$$

$$y = jk; \quad j = 0, 1, 2, 3, \dots$$

The points of intersection of these families of lines are called mesh points, lattice points or grid points.

Then:

$$u_x = \frac{u_{i+1,j} - u_{i,j}}{k} + 0(h) \text{ Forward differenc}$$

$$u_x = \frac{u_{i,j} - u_{i-1,j}}{h} + 0(h) \text{ Backward differenc}$$

$$u_x = \frac{u_{i+1,j} - u_{i-1,j}}{2h} + 0(h^2) \text{ Central differenc}$$

and;

$$u_{xx} = \frac{u_{i-1,j} - 2u_{i,j} + u_{i+1,j}}{h^2} + 0(h^2)$$

Where  $u_{i,j} = u(ih, jk) = u(x, y)$

*Similarly we have the approximations:*

$$u_y = \frac{u_{i,j+1} - u_{i,j}}{k} + O(k)$$

*Forward differenc*

$$u_y = \frac{u_{i,j} - u_{i,j-1}}{k} + O(k)$$

*Backward differenc*

$$u_y = \frac{u_{i,j+1} - u_{i,j-1}}{2k} + O(k^2)$$

*Central differenc*

*and;*

$$u_{yy} = \frac{u_{i,j-1} - 2u_{i,j} + u_{i,j+1}}{k^2} + O(k^2)$$

We can now obtain the finite-difference analogues of partial differential equations by replacing the derivatives in any equation by their corresponding difference approximation given above.

*Thus, the Laplace equation in two dimension, namely,*

$$u_{xx} + u_{yy} = 0$$

*has its finite-difference analogue;*

$$\frac{1}{h^2} [u_{i+1,j} - 2u_{i,j} + u_{i-1,j}] + \frac{1}{k^2} [u_{i,j+1} - 2u_{i,j} + u_{i,j-1}] = 0$$

*If  $h = k$ , this gives,*

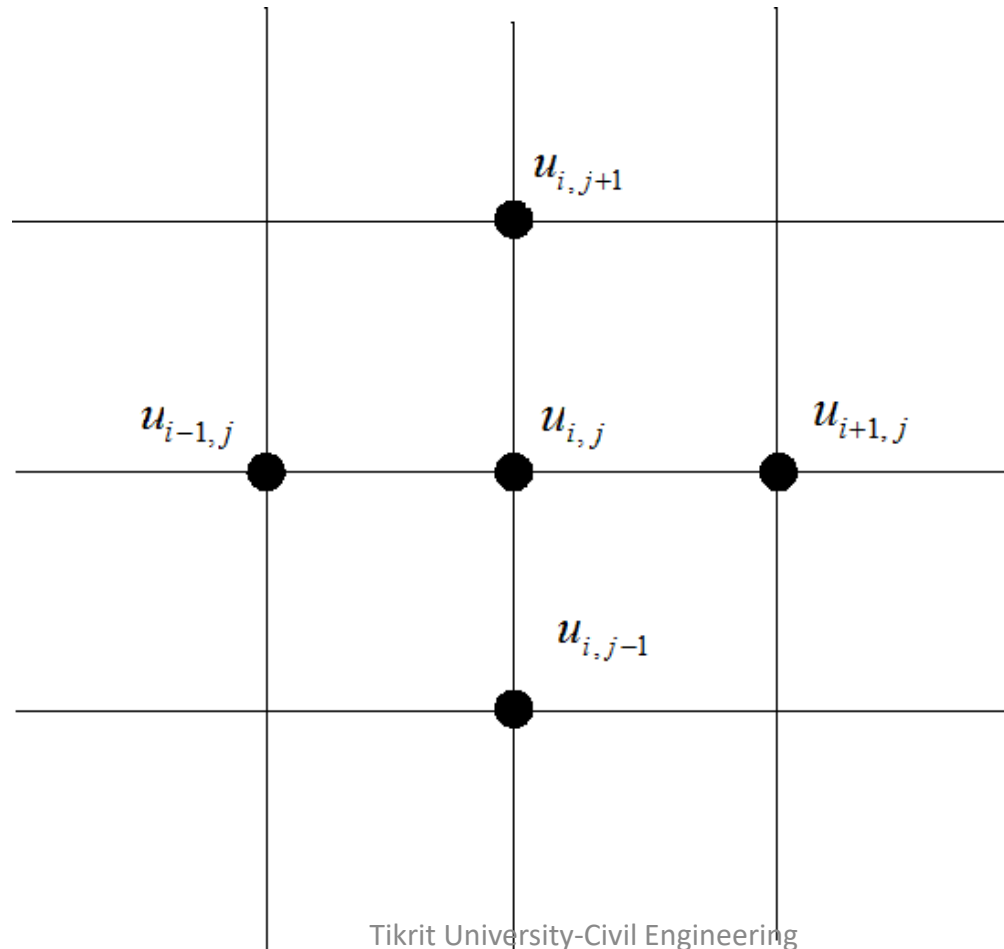
$$u_{i,j} = \frac{1}{4} [u_{i+1,j} + u_{i-1,j} + u_{i,j+1} + u_{i,j-1}] \quad (a)$$

*Which shows that the value of  $u$  at any point is the mean of its values at the four neighbouring points. This is called the standard five-point formula.*

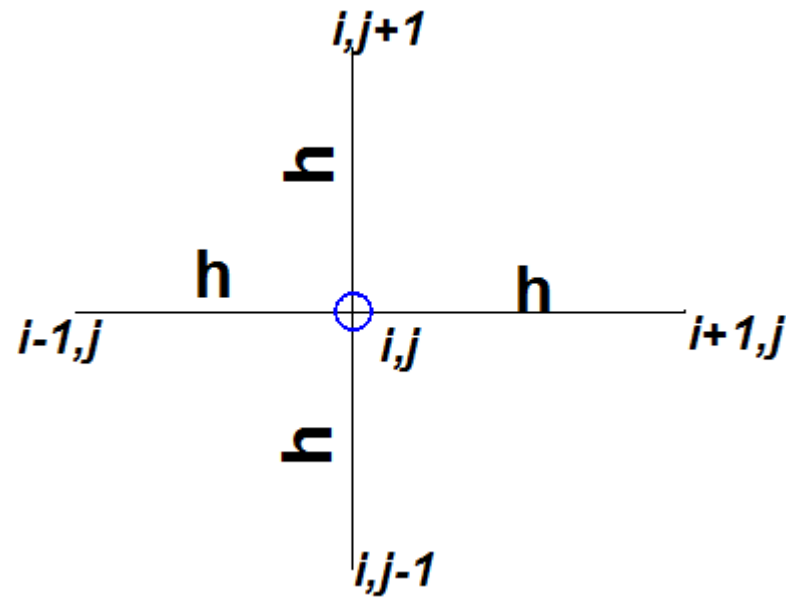


The standard five-point formula is written:

$$u_{i+1,j} + u_{i-1,j} + u_{i,j+1} + u_{i,j-1} - 4u_{i,j} = 0$$



Dirichlet  
problem



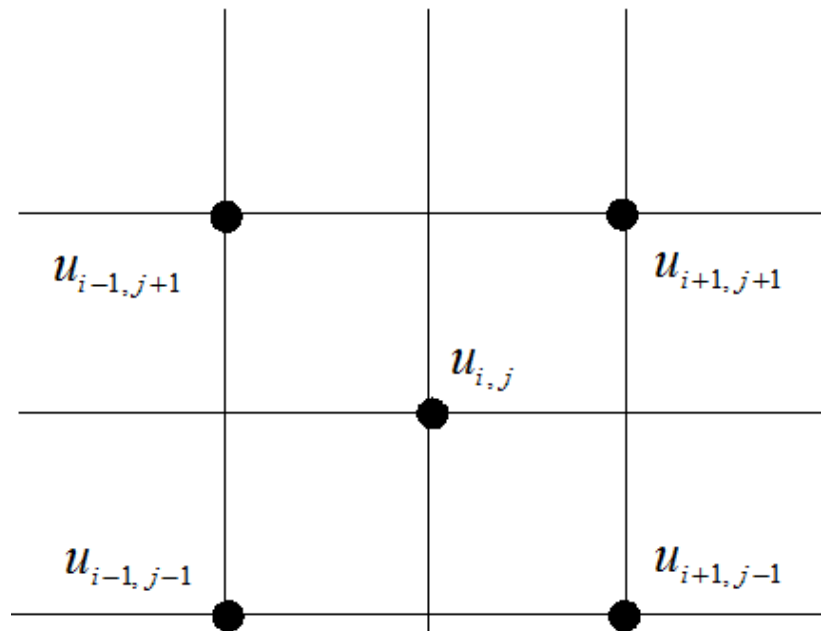
$$\begin{Bmatrix} 1 \\ 1 & -4 & 1 \\ 1 \end{Bmatrix}$$

*Also, instead of formula (a), we may use the following formula:*

$$u_{i,j} = \frac{1}{4} [u_{i-1,j-1} + u_{i+1,j-1} + u_{i+1,j+1} + u_{i-1,j+1}]$$

*Which uses the function values at the diagonal points.*

*And is therefore called the diagonal five-point formula.*

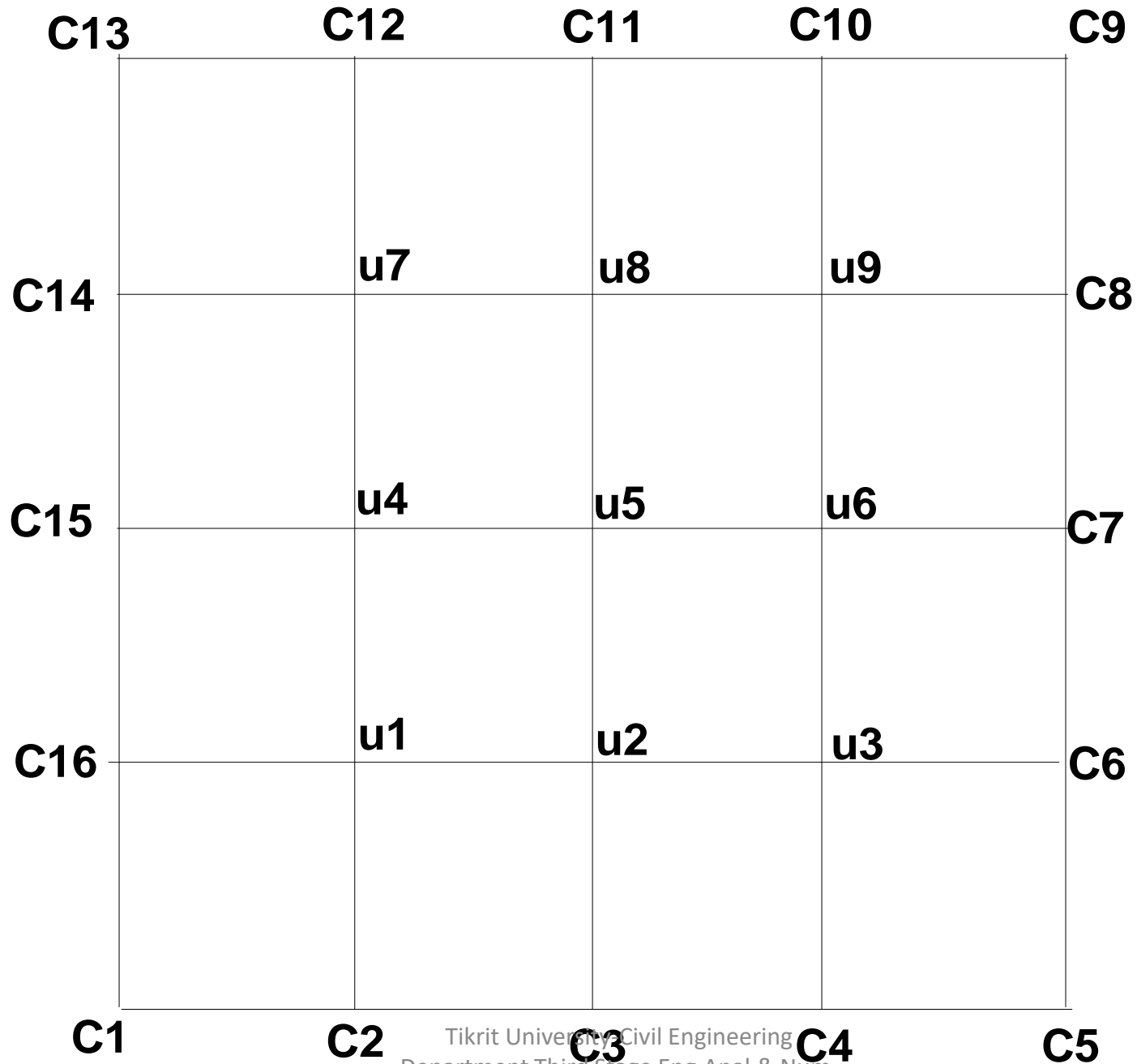


## Laplace's equation:

We wish to solve Laplace's equation:

$$u_{xx} + u_{yy} = 0$$

in a boundary region  $R$  with boundary  $C$ . As in Dirichlet's problem, let the value of  $u$  be specified every where on  $C$ . For simplicity, let  $R$  be a square region so that it can be divided into a network of small squares of side  $h$ . Let the values of  $u(x, y)$  on the boundary  $C$  be given by  $C_i$  and let the interior mesh points and the boundary points be as in the figure below:



*We first use the diagonal five – point formula and compute,  $u_5, u_7, u_9, u_1$  and  $u_3$  in this order, Thus we obtain,*

$$u_5 = \frac{1}{4}[C_1 + C_5 + C_9 + C_{13}];$$

$$u_7 = \frac{1}{4}[C_{15} + u_5 + C_{11} + C_{13}];$$

$$u_9 = \frac{1}{4}[u_5 + C_7 + C_9 + C_{11}];$$

$$u_1 = \frac{1}{4}[C_1 + C_3 + u_5 + C_{15}] \text{ and}$$

$$u_3 = \frac{1}{4}[C_3 + C_5 + C_7 + u_5]$$

*We then compute, in this order, the remaining quantities such  $u_8, u_4, u_6$  and  $u_2$  by the standard five-point formula. Thus we have,*

$$u_8 = \frac{1}{4}[u_5 + u_9 + C_{11} + C_7];$$

$$u_4 = \frac{1}{4}[u_1 + u_5 + u_7 + C_{15}];$$

$$u_6 = \frac{1}{4}[u_5 + u_3 + C_7 + u_9] \text{ and}$$

$$u_2 = \frac{1}{4}[C_3 + u_3 + u_5 + u_1]$$

*When once all the  $u_i (i = 1, 2, 3, \dots, 9)$  are computed, their accuracy can be improved by any of the iterative methods described below:*

*(1) Jacobi's method :*

*Let  $u_{i,j}^{(n)}$  denote the  $n^{th}$  iterative value of  $u_{i,j}$ .*

*An iterative procedure to solve the Eq.(a) is :*

$$u_{i,j}^{(n+1)} = \frac{1}{4} [u_{i-1,j}^{(n)} + u_{i+1,j}^{(n)} + u_{i,j-1}^{(n)} + u_{i,j+1}^{(n)}]$$

*for the interior mesh points. This is called the point Jacobi method.*



## (2) Gauss – Seidal method:

*The method uses the latest iterative values available and scans the mesh points systematically from left to right along successive rows.*

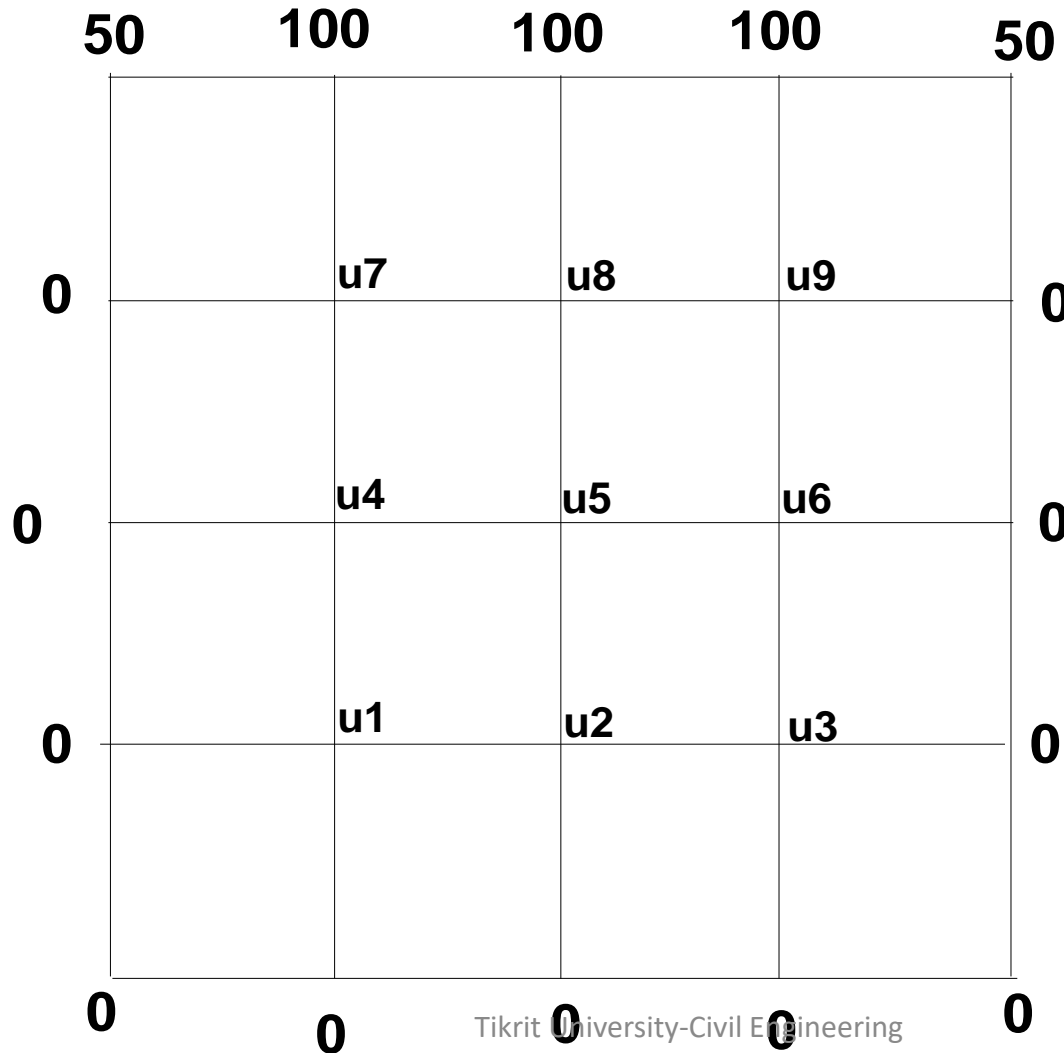
*The iterative formula is:*

$$u_{i,j}^{(n+1)} = \frac{1}{4} [u_{i-1,j}^{(n+1)} + u_{i+1,j}^{(n)} + u_{i,j-1}^{(n+1)} + u_{i,j+1}^{(n)}]$$

*It can be shown that the Gauss – Seidal scheme converges twice as fast as the Jacobi scheme.*

## Example:

Solve Laplace equation for the figure given below:



### Solution:

*We first compute the quantities  $u_5, u_7, u_9, u_1$  and  $u_3$  by using the diagonal five-point formula:*

$$u_5^{(1)} = \frac{1}{4}[0 + 0 + 50 + 50] = 25$$

$$u_7^{(1)} = \frac{1}{4}[0 + 25 + 100 + 50] = 43.75$$

$$u_9^{(1)} = \frac{1}{4}[25 + 0 + 50 + 100] = 43.75$$

$$u_1^{(1)} = \frac{1}{4}[0 + 0 + 25 + 0] = 6.25$$

$$u_3^{(1)} = \frac{1}{4}[0 + 0 + 0 + 25] = 6.25$$

*We now compute;  $u_8, u_4, u_6$  and  $u_2$  successively by using the standard five-point formula:*

$$u_8^{(1)} = \frac{1}{4}[25 + 43.75 + 100 + 43.75] = 53.125$$

$$u_4^{(1)} = \frac{1}{4}[0 + 6.25 + 25 + 43.75] = 18.75$$

$$u_6^{(1)} = \frac{1}{4}[25 + 6.25 + 0 + 43.75] = 18.75$$

$$u_2^{(1)} = \frac{1}{4}[6.25 + 0 + 6.25 + 25] = 9.375$$

*We have thus obtained the first approximations of all the nine mesh points and we can now use one of the iterative formula, by using Gauss–Seidal formula:*

$$u_{i,j}^{(n+1)} = \frac{1}{4} [u_{i-1,j}^{(n+1)} + u_{i+1,j}^{(n)} + u_{i,j-1}^{(n+1)} + u_{i,j+1}^{(n)}]$$

n	$u_1$	$u_2$	$u_3$	$u_4$	$u_5$	$u_6$	$u_7$	$u_8$	$u_9$
1	6.25	9.38	6.25	18.75	25.00	18.75	43.75	53.13	43.75
2	7.03	9.57	7.08	18.94	25.10	18.98	43.02	52.97	42.99
3	7.13	9.83	7.20	18.81	25.15	18.84	42.94	52.77	42.90
4	7.16	9.88	7.18	18.81	25.08	18.79	42.89	52.72	42.88
5	7.17	9.86	7.16	18.78	25.04	18.77	42.88	52.70	42.87

$$u_{1,1}^{(2)} = \frac{1}{4}[u_{0,1}^{(2)} + u_{2,1}^{(1)} + u_{1,0}^{(2)} + u_{1,2}^{(1)}]$$

$$u_1 = \frac{1}{4}[0 + 9.38 + 0 + 18.75] = 7.03$$

$$u_{2,1}^{(2)} = \frac{1}{4}[u_{1,1}^{(2)} + u_{3,1}^{(1)} + u_{2,0}^{(2)} + u_{2,2}^{(1)}]$$

$$u_2 = \frac{1}{4}[7.03 + 6.25 + 0 + 25] = 9.57$$